Search for Narrow Hγ Resonances in Proton-Proton Collisions at $\sqrt{s} = 13$ TeV

CMS Collaboration ; Canelli, Maria Florencia ; Kilminster, Benjamin ; Aarrestad, Thea K ; Brzhechko, Danyyl ; Caminada, Lea ; de Cosa, Annapaoloa ; Del Burgo, Riccardo ; Donato, Silvio ; Galloni, Camilla ; Hreus, Tomas ; Leontsinis, Stefanos ; Mikuni, Vinicius Massami ; Neutelings, Izaak ; Rauco, Giorgia ; Robmann, Peter ; Salerno, Daniel ; Schweiger, Korbinian ; Seitz, Claudia ; Takahashi, Yuta ; Wertz, Sebastien ; Zucchetta, Alberto ; et al

Abstract: A search for heavy, narrow resonances decaying to a Higgs boson and a photon (H) has been performed in proton-proton collision data at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 35.9 fb^1 collected with the CMS detector at the LHC in 2016. Events containing a photon and a Lorentz-boosted hadronically decaying Higgs boson reconstructed as a single, large-radius jet are considered, and the +jet invariant mass spectrum is analyzed for the presence of narrow resonances. To increase the sensitivity of the search, events are categorized depending on whether or not the large-radius jet can be identified as a result of the merging of two jets originating from b quarks. Results in both categories are found to agree with the predictions of the standard model. Upper limits on the production rate of H resonances are set as a function of their mass in the range of 720–3250 GeV, representing the most stringent constraints to date.

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Search for Narrow $H\gamma$ Resonances in Proton-Proton Collisions at $\sqrt{s} = 13$ TeV

A. M. Sirunyan *et al.*[*]
(CMS Collaboration)

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A search for heavy, narrow resonances decaying to a Higgs boson and a photon ($H\gamma$) has been performed in proton-proton collision data at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 35.9 fb^{-1} collected with the CMS detector at the LHC in 2016. Events containing a photon and a Lorentz-boosted hadronically decaying Higgs boson reconstructed as a single, large-radius jet are considered, and the $\gamma +$ jet invariant mass spectrum is analyzed for the presence of narrow resonances. To increase the sensitivity of the search, events are categorized depending on whether or not the large-radius jet can be identified as a result of the merging of two jets originating from b quarks. Results in both categories are found to agree with the predictions of the standard model. Upper limits on the production rate of $H\gamma$ resonances are set as a function of their mass in the range of 720–3250 GeV, representing the most stringent constraints to date.

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The CERN LHC has enabled the investigation of numerous theories beyond the standard model (SM), among which are those that predict the production of new heavy bosons that decay into a pair of SM bosons. These resonances are found in a large variety of models, from theories with an extended Higgs sector, such as the two Higgs doublet models [1], to models with low-scale gravity, e.g., the Randall-Sundrum model [2,3], and models with strong dynamics [4,5]. A number of searches for resonances decaying into the WW, WZ, ZZ, $W\gamma$, or $Z\gamma$ channels have been carried out at 7, 8, and 13 TeV center-of-mass energies [6–25] by the ATLAS and CMS Collaborations, and stringent limits on their production have been set. The discovery of a Higgs boson (H) [26–28] allowed for an extension of these searches to include the VH channels, where V stands for a W or Z boson. Negative results of these searches at 8 and 13 TeV have been reported by CMS [16,29–33] and ATLAS [34–37].

In this Letter, we extend the above program by reporting on a search for narrow resonances decaying into a Higgs boson and a photon ($H\gamma$). The existence of such resonances has been recently predicted [38], using a $Z' \to H\gamma$ decay at one-loop level as an example, which we chose as the signal benchmark in the present analysis. Here, Z' is a $U(1)$ spin-1 boson, similar to the SM Z boson. For the decay of a sufficiently massive Z' boson via this channel, the Higgs boson would be produced with a significant Lorentz boost. For the dominant Higgs boson decay mode into a pair of b quarks, recently established by ATLAS and CMS [39,40], the decay products will be reconstructed as a single large-radius jet with a characteristic substructure that can be exploited to distinguish those jets from jets originating from gluons or light-flavor quarks. The background to this process is dominated by SM $\gamma +$ jet production, and the application of "b tagging" techniques to identify jets originating from b quarks ("b quark jets" or "b jets") can reduce the backgrounds by nearly two orders of magnitude. The remaining backgrounds come from nonresonant $b\bar{b}\gamma$ production, from $q\bar{q}\gamma$ production with the light-flavor quarks incorrectly identified as b quarks, and from multijet production with light-flavor jets mistagged as b quark jets in addition to a jet misreconstructed as a photon. Very recently, when the work presented here was nearing completion, ATLAS reported on a first search in this channel [41].

The analysis described in this Letter is performed with data corresponding to an integrated luminosity of 35.9 fb^{-1} recorded with the CMS detector at the LHC in proton-proton (pp) collisions at a center-of-mass energy of 13 TeV in 2016. We focus on Z' masses above ≈700 GeV where the Higgs boson is produced with a sufficient boost to be efficiently reconstructed as a single, large-radius jet. The analysis uses control samples in data to predict the shape of the background and to optimize the selection. The background is obtained via a fit to data with a functional form determined using these control samples and simulation, and a signal is searched for as a narrow resonance on top of a smoothly falling background. To increase the sensitivity of the search, events are categorized depending on whether or not the large-radius jet can be identified as a result of

[*]Full author list given at the end of the Letter.

the merging of two jets originating from b quarks. This categorization uses an advanced b tagging technique developed specifically for searches involving Lorentz-boosted Higgs bosons decaying into $b\bar{b}$ pairs [42].

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two end cap sections. Forward calorimeters extend the pseudorapidity (η) coverage provided by the barrel and end cap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid. Events of interest are selected using a two-tiered trigger system [43]. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a time interval of less than 4 μs. The second level, referred to as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to less than 1 kHz before data storage. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [44].

Online, events are selected using a logical OR of two triggers: one requiring a photon candidate with a transverse momentum (p_T) greater than 175 GeV and the other requiring a photon with $p_T > 165$ GeV, which in addition must have the ratio of the energy deposits in the HCAL and ECAL below 0.1. Offline, events are reconstructed using a particle-flow (PF) algorithm [45], which aims to reconstruct and identify each individual particle (charged or neutral hadron, photon, electron, or muon) in an event, with an optimized combination of information from the various elements of the CMS detector.

For each event, hadronic jets are clustered from PF candidates using the anti-k_T algorithm [46,47], with a distance parameter of 0.8 ("large-radius jet," or J), as implemented in FastJet [47]. The jet momentum is determined as the vectorial sum of all particle momenta in the jet, and is found from simulation to be within 5% to 10% of the true momentum over the whole p_T spectrum and detector acceptance. Additional pp interactions within the same or nearby bunch crossings (pileup) can contribute additional tracks and calorimetric energy depositions to the jet momentum. To mitigate this effect, the pileup-per-particle identification algorithm [48] is applied to PF candidates prior to jet clustering. Jet energy corrections are derived from both simulation and data to make the measured response of jets equal to that of particle-level jets on average [49]. The jet energy resolution amounts typically to 8% at 100 GeV, and 4% at 1000 GeV.

Photons are identified with a multivariate analysis (MVA) classifier [50] based on shower shape variables, isolation sums computed from PF candidates in a cone of radius $\Delta R = 0.3$ in the η-ϕ plane, centered on the photon candidate, and variables that account for the dependencies of the shower shape and isolation variables on pileup. In addition, a conversion-safe electron veto [50] is applied. Photon candidates are required to satisfy an MVA working point requirement that corresponds to a typical photon reconstruction and identification efficiency of 90% in the photon p_T range used in the analysis.

At the preselection level, events are required to contain at least one photon candidate with $p_T^\gamma > 200$ GeV and $|\eta^\gamma| < 2.4$, excluding the transition region of $1.44 < |\eta| < 1.57$ between the ECAL barrel and end cap, where the reconstruction is not optimal. At least one large-radius jet with $p_T^J > 250$ GeV and $|\eta^J| < 2.6$ is also required. The photon and jet must be separated by $\Delta R(\gamma, J) > 1.1$, so as to ensure no overlap between the jet and the photon isolation cone. The triggers are found to be more than 98% efficient for the signal events passing the preselection.

In addition, to get a more precise estimate of the Higgs boson jet invariant mass, a jet grooming algorithm known as "soft drop'" (SD) [51] is applied. The groomed jet mass (m_J^{SD}) is computed from the sum of the four-momenta of the remaining jet constituents, which are corrected with the same factor as has already been used in the generic jet reconstruction described above. The typical mass resolution for the Higgs boson jet is 10% [52].

To take advantage of the dominant Higgs boson decay mode, $H \to b\bar{b}$, we further classify the events according to the output of the dedicated double b tagging (DBT) algorithm [42], which attempts to identify a two-prong substructure within a large-radius jet, as well as the likelihood that the two subjets originate from the b quarks. For the latter, several associated variables are used in an MVA tagging algorithm. These variables include the significance of the impact parameters of the tracks relative to the primary vertex, the number and masses of secondary vertices, and a variable that characterizes the system of two secondary vertices, taking into account kinematic properties that help to distinguish $b\bar{b}$ pairs produced in massive-particle decays from those originating from gluon splitting. We define two event categories using the "tight" working point of the DBT algorithm [42], which corresponds to a 36 (20)% Higgs boson jet tagging efficiency for a signal mass of 1000 (3000) GeV and an ≈1% QCD jet mistag probability in events passing the preselection. This working point was found to give the greatest signal sensitivity at lower signal masses, below ≈2000 GeV, as it provides powerful background rejection. The events with the leading jet passing the tight working point of the DBT algorithm are assigned to the "b-tagged" category, while the rest are classified as "untagged." The untagged category allows the search to maintain optimal sensitivity to

high-mass resonances (above ~2000 GeV), for which the background is small and at the same time the DBT efficiency deteriorates compared to that at low masses, as the tracks originating from the secondary vertices become more collimated and harder to resolve.

For each category, two regions in preselected data are defined: the search region (SR), in which the invariant mass of the leading jet is required to be $110 < m_J^{\mathrm{SD}} < 140$ GeV, centered on the nominal Higgs boson mass, and the sideband (SB) region, which requires $100 < m_J^{\mathrm{SD}} < 110$ GeV. The SB region is chosen because events within it reproduce the shape of the distributions of dominant backgrounds in the SR for the kinematic variables used in the analysis, as seen in data and also confirmed by the Monte Carlo (MC) simulation. Consequently, the SB region can be used to predict the background shape in the SR, and therefore to optimize the analysis without reliance on background simulation. The lower boundary of the SB region is chosen so as to avoid contamination from W and Z dijet decays, and the upper boundary is chosen far enough from the nominal Higgs boson mass to minimize possible signal contamination. Since the mass of a large-radius jet is loosely correlated with its p_T, we find that with the preselection kinematic requirements, an adequate description of the background shape in the SR is achieved via the SB region for $J\gamma$ masses above 720 GeV. Therefore, we search for narrow $H\gamma$ resonances in the 720–3250 GeV range, with the upper range determined by the point beyond which we expect to see no background events in data in either of the two categories.

Simulated Z' signal samples are generated based on the benchmark model of Ref. [38], as implemented in the MADGRAPH5_aMC@NLO 2.3.3 [53] MC generator. Signals are generated at leading order (LO) in the mass range of 650–3250 GeV, in steps of 100 GeV (650–850 GeV), 150 GeV (850–2050 GeV), or 400 GeV (2050–3250 GeV). The intrinsic width of the signal is chosen to be negligibly small compared to the experimental resolution. The signal shape is parametrized with the Crystal Ball function [54], which provides an adequate description in the entire mass range considered. For other values of the mass, the signal shape is smoothly interpolated between the Crystal Ball function parameter values derived for the simulated mass points [55].

While the background estimates come from data, we use simulated background samples to check that the kinematic variables used in the optimization are adequately described by the data in the SB region. The following simulated background samples are used: multijets, $\gamma +$ jets, $W +$ jets, and $Z/\gamma^* +$ jets, all generated at LO with MADGRAPH5_aMC@NLO. All the background samples are normalized using next-to-leading order cross sections and the integrated luminosity of the data sample.

The NNPDF3.0 [56] parton distribution functions (PDFs) are used for all simulated samples. The fragmentation and hadronization are described with PYTHIA version 8.205 [57] using the underlying event tune CUETP8M1 [58]. The CMS detector response is simulated using GEANT4 [59]. All simulated samples include effects of pileup by superimposing on hard scattering events simulated minimum bias collisions, with their multiplicity matching that observed in data.

Scale factors are applied to simulated samples to remove discrepancies between various efficiencies in simulation compared to those in data. These scale factors range 0.85–0.91 [42] for the DBT efficiency, and ≈0.99 [50] for the photon identification efficiency.

Further requirements on the kinematic variables of selected events are applied in order to ensure optimal signal sensitivity. The background estimate for the optimization studies comes from the SB region, normalized to the overall number of events in the SR. The optimization is performed separately for the b-tagged and untagged categories. Both the signal significance and the signal limit optimization were checked and in most of the cases the optimal points are the same. Several kinematic variables were considered, including the N-subjettiness variables [60], characterizing the substructure of the jet. As a result, the following additional requirements were chosen on top of the preselection for both b-tagged and untagged categories: the leading jet must have $|\eta| < 2.2$, ensuring that the core of the jet is within the tracker coverage, the photon must be found in the ECAL barrel ($|\eta| < 1.44$), and the ratio of the photon p_T to the mass of the $J\gamma$ system, $p_T^{\gamma}/m_{J\gamma}$, must exceed 0.35. The last two requirements suppress photons from the $\gamma +$ jet background, which tend to be more forward than the signal photons.

Shown in Fig. 1 are the products of signal acceptance and efficiency versus generated signal mass for the two analysis categories, evaluated for each of the simulated signal samples, and a fit function used to interpolate between the generated mass points. In the b-tagged category, the product of the overall acceptance, and the reconstruction, trigger, and full selection efficiency for signal events increases from about 3% at low signal masses to a peak

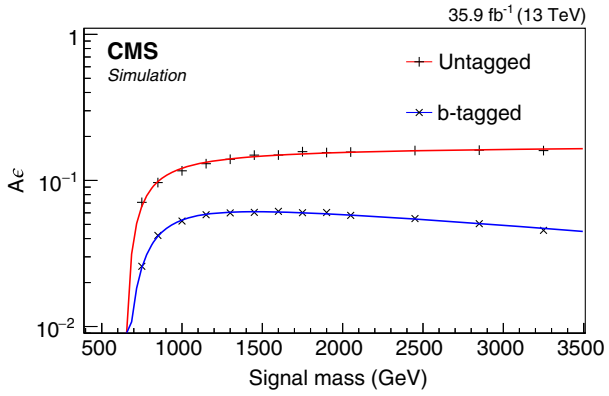

FIG. 1. Signal acceptance times efficiency ($A\epsilon$) after the final selection, shown for the b-tagged and untagged categories.

of 6% near a signal invariant mass of 1500 GeV, and decreases thereafter to about 5% for high signal masses. The observed behavior at high masses is due to the degradation in the DBT efficiency. For the untagged category, the corresponding product of acceptance and efficiency increases from about 7% at low signal masses to around 16% for high signal masses. The main factors that impact the acceptance are the 57% SM Higgs boson branching fraction to $b\bar{b}$ and the fact that about 35% of the signal large-radius jets fail the $110 < m_J^{\mathrm{SD}} < 140$ GeV SR requirement.

After the final selection, the background shape in each category is modeled by fitting a smooth, monotonically falling function to the $J\gamma$ invariant mass spectrum in the SR. A variety of functional forms are considered for the background fit, based on the SB region data and on simulated background samples. For every function, a goodness-of-fit (GOF) test known as the method of saturated models [61] is performed in the SR. The nominal background fit function is then chosen as the one with the best GOF with the minimal number of parameters. The selection of the nominal fit function is performed independently for the b-tagged and untagged categories. In both categories, the following function was found to give the best GOF to the $J\gamma$ invariant mass spectrum in the SR: $dN/dm = p_0(m/\sqrt{s})^{p_1+p_2\log(m/\sqrt{s})}$, where p_i, $i = 0, 1, 2$ are the free parameters of the fit.

In order to prove that no systematic bias arises because of the choice of the background fit function, a number of tests are performed. An alternative fit function that performed well in the GOF test is used to generate a large number of $J\gamma$ invariant mass spectra, with or without signal injection. The spectra are then fit to the sum of the chosen background template and a signal with the mass and normalization allowed to float. The signal significance is extracted from each fit and the distributions of the pull of the signal yield are constructed, where the pull is defined as the difference between the injected and extracted signal normalizations, divided by the statistical uncertainty in the extracted signal normalization from the fit. We observe that the distributions of the pulls are consistent with a Gaussian function with zero mean and a standard deviation of unity, and thus conclude that any systematic bias from the background fitting procedure is negligible compared to the statistical uncertainties in the fit. We therefore use the latter as the only uncertainties associated with the background estimate.

Several systematic uncertainties are taken into account in the signal extraction procedure. These uncertainties stem from effects that may lead to an imperfect estimate of the signal rate and shape, including experimental uncertainties in the integrated luminosity (2.5%) [62]), jet energy scale and resolution (2.0%) [49,63], photon energy scale and resolution (0.1–2.3%, depending on the $J\gamma$ mass [22]), pileup (1.0%), groomed jet mass scale (5.0%), and various identification efficiencies (4.0%, dominated by the DBT efficiency uncertainty [42]). We also include an uncertainty in signal acceptance due to the PDF choice (2.0%), based on the PDF4LHC recommendations [64] using the NNPDF3.0 replicas [56]. Since the correction for the trigger inefficiency in data never exceeds 2%, the uncertainty due to this

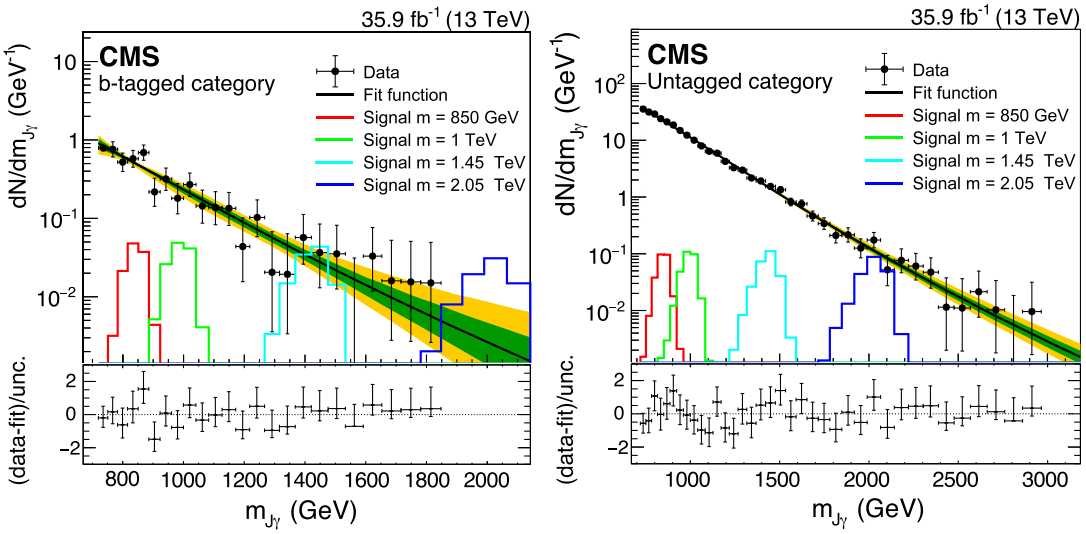

FIG. 2. The observed $J\gamma$ invariant mass spectra in the signal region, shown along with the background fit and a few selected signals, for the b-tagged (left) and untagged (right) categories. Signal samples are plotted with arbitrary normalizations and are shown for illustration purposes. The green and yellow bands correspond to the 1 and 2 standard deviation uncertainties in the background-only fit. For bins with a low number of data entries, the error bars correspond to the Garwood confidence intervals [65]. Shown in the lower panels are the differences between the number of events in data and the nominal background prediction from the fit, divided by the combined statistical uncertainty in the data and the background fit. The error bars correspond to the statistical uncertainty in the data alone.

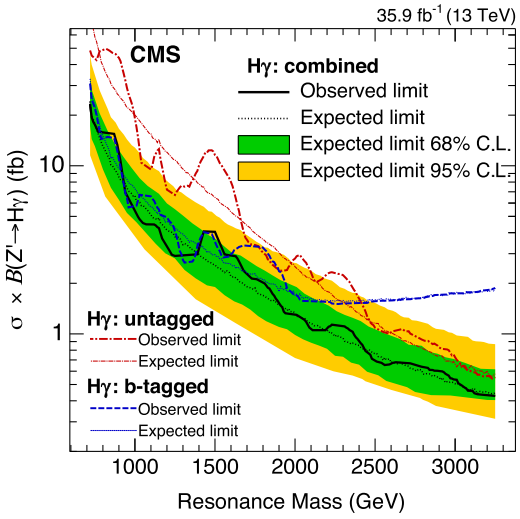

FIG. 3. Upper limits at 95% C.L. on the product of the signal cross section and the branching fraction to $H\gamma$ for the b-tagged and untagged categories, as well as the statistical combination of the two categories.

correction is always much smaller than the statistical uncertainty of data and therefore has been ignored.

The $J\gamma$ invariant mass spectra for both b-tagged and untagged categories, together with the background-only fit, as well as expected signal shapes for several signal masses, are shown in Fig. 2. The background fit function is shown with the 68% and 95% confidence level (C.L.) uncertainty bands obtained from the fit. Results in both categories are found to agree with the background-only hypothesis and do not exhibit any significant resonance-like structures. We set upper limits on the production cross section of narrow spin-1 resonances using the modified frequentist CL_s criterion [66–68], with a likelihood ratio used as a test statistic, and uncertainties incorporated as nuisance parameters with log-normal priors.

Shown in Fig. 3 are 95% C.L. upper limits on the product of the signal cross section and the branching fraction to $H\gamma$ for the b-tagged category, the untagged category, and the statistical combination of the two categories. These limits are the most stringent to date in the entire mass range studied and are the only available limits for masses below 1000 GeV and above 3000 GeV. The significant improvement in the sensitivity compared to the very recent ATLAS limits [41] in the 1000–3000 GeV range results from the application of the more efficient DBT algorithm, at low masses, and from the use of the untagged event category, at high masses.

In summary, a search for heavy, narrow resonances decaying to a Higgs boson and a photon ($H\gamma$) has been performed in proton-proton collision data at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 35.9 fb^{-1} collected with the CMS detector at the LHC in 2016. Events in which a photon and a Lorentz-boosted Higgs boson that decays hadronically and is reconstructed as a single, large-radius jet are considered, and the γ + jet invariant mass spectrum is analyzed for the presence of narrow resonances. To increase the sensitivity of the search, events are categorized depending on whether the large-radius jet can be identified as a result of the merging of two jets originating from b quarks. The backgrounds, dominated by standard model γ + jet production, are estimated directly from data, without reliance on simulation. Results in both categories are found to agree with the predictions of the standard model. Upper limits on the production cross section of $H\gamma$ resonances ranging from 25 to 0.4 fb are set as a function of the resonance mass in the range of 720–3250 GeV. These are the most stringent constraints on narrow, spin-1 $H\gamma$ resonances to date in the entire mass range, and the first limits available below 1000 GeV and above 3000 GeV.

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Molnar,[50] Z. Szillasi,[50] P. Raics,[51] Z. L. Trocsanyi,[51] B. Ujvari,[51] S. Choudhury,[52] J. R. Komaragiri,[52] P. C. Tiwari,[52] S. Bahinipati,[53,x] C. Kar,[53] P. Mal,[53] K. Mandal,[53] A. Nayak,[53,y] D. K. Sahoo,[53,x] S. K. Swain,[53] S. Bansal,[54] S. B. Beri,[54] V. Bhatnagar,[54] S. Chauhan,[54] R. Chawla,[54] N. Dhingra,[54] R. Gupta,[54] A. Kaur,[54] M. Kaur,[54] S. Kaur,[54] R. Kumar,[54] P. Kumari,[54] M. Lohan,[54] A. Mehta,[54] K. Sandeep,[54] S. Sharma,[54] J. B. Singh,[54] A. K. Virdi,[54] G. Walia,[54] A. Bhardwaj,[55] B. C. Choudhary,[55] R. B. Garg,[55] M. Gola,[55] S. Keshri,[55] Ashok Kumar,[55] S. Malhotra,[55] M. Naimuddin,[55] P. Priyanka,[55] K. Ranjan,[55] Aashaq Shah,[55] R. Sharma,[55] R. Bhardwaj,[56,z] M. Bharti,[56,z] R. Bhattacharya,[56] S. Bhattacharya,[56] U. Bhawandeep,[56,z] D. Bhowmik,[56] S. Dey,[56] S. Dutt,[56,z] S. Dutta,[56] S. Ghosh,[56] K. Mondal,[56] S. Nandan,[56] A. Purohit,[56] P. K. Rout,[56] A. Roy,[56] S. Roy Chowdhury,[56] G. Saha,[56] S. Sarkar,[56] M. Sharan,[56] B. Singh,[56,z] S. Thakur,[56,z] P. K. Behera,[57] R. Chudasama,[58] D. Dutta,[58] V. Jha,[58] V. Kumar,[58] P. K. Netrakanti,[58] L. M. Pant,[58] P. Shukla,[58] T. Aziz,[59] M. A. Bhat,[59] S. Dugad,[59] G. B. Mohanty,[59] N. Sur,[59] B. Sutar,[59] Ravindra Kumar Verma,[59] S. Banerjee,[60] S. Bhattacharya,[60] S. Chatterjee,[60] P. Das,[60] M. Guchait,[60] Sa. Jain,[60] S. Karmakar,[60] S. Kumar,[60] M. Maity,[60,aa] G. Majumder,[60] K. Mazumdar,[60] N. Sahoo,[60] T. Sarkar,[60,aa] S. Chauhan,[61] S. Dube,[61] V. Hegde,[61] A. Kapoor,[61] K. Kothekar,[61] S. Pandey,[61] A. Rane,[61] S. Sharma,[61] S. Chenarani,[62,bb] E. Eskandari Tadavani,[62] S. M. Etesami,[62,bb] M. Khakzad,[62] M. Mohammadi Najafabadi,[62] M. Naseri,[62] F. Rezaei Hosseinabadi,[62] B. Safarzadeh,[62,cc] M. Zeinali,[62] M. Felcini,[63] M. Grunewald,[63] M. Abbrescia,[64a,64b] C. Calabria,[64a,64b] A. Colaleo,[64a] D. Creanza,[64a,64c] L. Cristella,[64a,64b] N. De Filippis,[64a,64c] M. De Palma,[64a,64b] A. Di Florio,[64a,64b] F. Errico,[64a,64b] L. Fiore,[64a] A. Gelmi,[64a,64b] G. Iaselli,[64a,64c] M. Ince,[64a,64b] S. Lezki,[64a,64b] G. Maggi,[64a,64c] M. Maggi,[64a] G. Miniello,[64a,64b] S. My,[64a,64b] S. Nuzzo,[64a,64b] A. Pompili,[64a,64b] G. Pugliese,[64a,64c] R. Radogna,[64a] A. Ranieri,[64a] G. Selvaggi,[64a,64b] A. Sharma,[64a] L. Silvestris,[64a] R. Venditti,[64a] P. Verwilligen,[64a] G. Zito,[64a] G. Abbiendi,[65a] C. Battilana,[65a,65b] D. Bonacorsi,[65a,65b] L. Borgonovi,[65a,65b] S. Braibant-Giacomelli,[65a,65b] R. Campanini,[65a,65b] P. Capiluppi,[65a,65b] A. Castro,[65a,65b] F. R. Cavallo,[65a] S. S. Chhibra,[65a,65b] C. Ciocca,[65a] G. Codispoti,[65a,65b] M. Cuffiani,[65a,65b] G. M. Dallavalle,[65a] F. Fabbri,[65a] A. Fanfani,[65a,65b] E. Fontanesi,[65a] P. Giacomelli,[65a] C. Grandi,[65a] L. Guiducci,[65a,65b] F. Iemmi,[65a,65b] S. Lo Meo,[65a] S. Marcellini,[65a] G. Masetti,[65a] A. Montanari,[65a] F. L. Navarria,[65a,65b] A. Perrotta,[65a] F. Primavera,[65a,65b,q] T. Rovelli,[65a,65b] G. P. Siroli,[65a,65b] N. Tosi,[65a] S. Albergo,[66a,66b] A. Di Mattia,[66a] R. Potenza,[66a,66b] A. Tricomi,[66a,66b] C. Tuve,[66a,66b] G. Barbagli,[67a] K. Chatterjee,[67a,67b] V. Ciulli,[67a,67b] C. Civinini,[67a] R. D'Alessandro,[67a,67b] E. Focardi,[67a,67b] G. Latino,[67a] P. Lenzi,[67a,67b] M. Meschini,[67a] S. Paoletti,[67a] L. Russo,[67a,dd] G. Sguazzoni,[67a] D. Strom,[67a] L. Viliani,[67a] L. Benussi,[68] S. Bianco,[68] F. Fabbri,[68] D. Piccolo,[68] F. Ferro,[69a] F. Ravera,[69a,69b] E. Robutti,[69a] S. Tosi,[69a,69b] A. Benaglia,[70a] A. Beschi,[70a,70b] F. Brivio,[70a,70b] V. Ciriolo,[70a,70b,q] S. Di Guida,[70a,70b,q] M. E. Dinardo,[70a,70b] S. Fiorendi,[70a,70b] S. Gennai,[70a] A. Ghezzi,[70a,70b] P. Govoni,[70a,70b] M. Malberti,[70a,70b] S. Malvezzi,[70a] A. Massironi,[70a,70b] D. Menasce,[70a] F. Monti,[70a] L. Moroni,[70a] M. Paganoni,[70a,70b] D. Pedrini,[70a] S. Ragazzi,[70a,70b] T. Tabarelli de Fatis,[70a,70b] D. Zuolo,[70a,70b] S. Buontempo,[71a] N. Cavallo,[71a,71c] A. De Iorio,[71a,71b] A. Di Crescenzo,[71a,71b] F. Fabozzi,[71a,71c] F. Fienga,[71a] G. Galati,[71a] A. O. M. Iorio,[71a,71b] W. A. Khan,[71a] L. Lista,[71a] S. Meola,[71a,71d,q] P. Paolucci,[71a,q] C. Sciacca,[71a,71b] E. Voevodina,[71a,71b] P. Azzi,[72a] N. Bacchetta,[72a] D. Bisello,[72a,72b] A. Boletti,[72a,72b] A. Bragagnolo,[72a] R. Carlin,[72a,72b] P. Checchia,[72a] M. Dall'Osso,[72a,72b] P. De Castro Manzano,[72a] T. Dorigo,[72a] U. Dosselli,[72a] U. Gasparini,[72a,72b] S. Y. Hoh,[72a] S. Lacaprara,[72a] P. Lujan,[72a] M. Margoni,[72a,72b] A. T. Meneguzzo,[72a,72b] M. Passaseo,[72a] J. Pazzini,[72a,72b] N. Pozzobon,[72a,72b]

P. Ronchese,[72a,72b] R. Rossin,[72a,72b] F. Simonetto,[72a,72b] A. Tiko,[72a] E. Torassa,[72a] M. Zanetti,[72a,72b] P. Zotto,[72a,72b] G. Zumerle,[72a,72b] A. Braghieri,[73a] A. Magnani,[73a] P. Montagna,[73a,73b] S. P. Ratti,[73a,73b] V. Re,[73a] M. Ressegotti,[73a,73b] C. Riccardi,[73a,73b] P. Salvini,[73a] I. Vai,[73a,73b] P. Vitulo,[73a,73b] M. Biasini,[74a,74b] G. M. Bilei,[74a] C. Cecchi,[74a,74b] D. Ciangottini,[74a,74b] L. Fanò,[74a,74b] P. Lariccia,[74a,74b] R. Leonardi,[74a,74b] E. Manoni,[74a] G. Mantovani,[74a,74b] V. Mariani,[74a,74b] M. Menichelli,[74a] A. Rossi,[74a,74b] A. Santocchia,[74a,74b] D. Spiga,[74a] K. Androsov,[75a] P. Azzurri,[75a] G. Bagliesi,[75a] L. Bianchini,[75a] T. Boccali,[75a] L. Borrello,[75a] R. Castaldi,[75a] M. A. Ciocci,[75a,75b] R. Dell'Orso,[75a] G. Fedi,[75a] F. Fiori,[75a,75c] L. Giannini,[75a,75c] A. Giassi,[75a] M. T. Grippo,[75a] F. Ligabue,[75a,75c] E. Manca,[75a,75c] G. Mandorli,[75a,75c] A. Messineo,[75a,75b] F. Palla,[75a] A. Rizzi,[75a,75b] P. Spagnolo,[75a] R. Tenchini,[75a] G. Tonelli,[75a,75b] A. Venturi,[75a] P. G. Verdini,[75a] L. Barone,[76a,76b] F. Cavallari,[76a] M. Cipriani,[76a,76b] D. Del Re,[76a,76b] E. Di Marco,[76a,76b] M. Diemoz,[76a] S. Gelli,[76a,76b] E. Longo,[76a,76b] B. Marzocchi,[76a,76b] P. Meridiani,[76a] G. Organtini,[76a,76b] F. Pandolfi,[76a] R. Paramatti,[76a,76b] F. Preiato,[76a,76b] S. Rahatlou,[76a,76b] C. Rovelli,[76a] F. Santanastasio,[76a,76b] N. Amapane,[77a,77b] R. Arcidiacono,[77a,77c] S. Argiro,[77a,77b] M. Arneodo,[77a,77c] N. Bartosik,[77a] R. Bellan,[77a,77b] C. Biino,[77a] N. Cartiglia,[77a] F. Cenna,[77a,77b] S. Cometti,[77a] M. Costa,[77a,77b] R. Covarelli,[77a,77b] N. Demaria,[77a] B. Kiani,[77a,77b] C. Mariotti,[77a] S. Maselli,[77a] E. Migliore,[77a,77b] V. Monaco,[77a,77b] E. Monteil,[77a,77b] M. Monteno,[77a] M. M. Obertino,[77a,77b] L. Pacher,[77a,77b] N. Pastrone,[77a] M. Pelliccioni,[77a] G. L. Pinna Angioni,[77a,77b] A. Romero,[77a,77b] M. Ruspa,[77a,77c] R. Sacchi,[77a,77b] K. Shchelina,[77a,77b] V. Sola,[77a] A. Solano,[77a,77b] D. Soldi,[77a,77b] A. Staiano,[77a] S. Belforte,[78a] V. Candelise,[78a,78b] M. Casarsa,[78a] F. Cossutti,[78a] A. Da Rold,[78a,78b] G. Della Ricca,[78a,78b] F. Vazzoler,[78a,78b] A. Zanetti,[78a] D. H. Kim,[79] G. N. Kim,[79] M. S. Kim,[79] J. Lee,[79] S. Lee,[79] S. W. Lee,[79] C. S. Moon,[79] Y. D. Oh,[79] S. I. Pak,[79] S. Sekmen,[79] D. C. Son,[79] Y. C. Yang,[79] H. Kim,[80] D. H. Moon,[80] G. Oh,[80] B. Francois,[81] J. Goh,[81,ee] T. J. Kim,[81] S. Cho,[82] S. Choi,[82] Y. Go,[82] D. Gyun,[82] S. Ha,[82] B. Hong,[82] Y. Jo,[82] K. Lee,[82] K. S. Lee,[82] S. Lee,[82] J. Lim,[82] S. K. Park,[82] Y. Roh,[82] H. S. Kim,[83] J. Almond,[84] J. Kim,[84] J. S. Kim,[84] H. Lee,[84] K. Lee,[84] K. Nam,[84] S. B. Oh,[84] B. C. Radburn-Smith,[84] S. h. Seo,[84] U. K. Yang,[84] H. D. Yoo,[84] G. B. Yu,[84] D. Jeon,[85] H. Kim,[85] J. H. Kim,[85] J. S. H. Lee,[85] I. C. Park,[85] Y. Choi,[86] C. Hwang,[86] J. Lee,[86] I. Yu,[86] V. Dudenas,[87] A. Juodagalvis,[87] J. Vaitkus,[87] I. Ahmed,[88] Z. A. Ibrahim,[88] M. A. B. Md Ali,[88,ff] F. Mohamad Idris,[88,gg] W. A. T. Wan Abdullah,[88] M. N. Yusli,[88] Z. Zolkapli,[88] J. F. Benitez,[89] A. Castaneda Hernandez,[89] J. A. Murillo Quijada,[89] H. Castilla-Valdez,[90] E. De La Cruz-Burelo,[90] M. C. Duran-Osuna,[90] I. Heredia-De La Cruz,[90,hh] R. Lopez-Fernandez,[90] J. Mejia Guisao,[90] R. I. Rabadan-Trejo,[90] M. Ramirez-Garcia,[90] G. Ramirez-Sanchez,[90] R. Reyes-Almanza,[90] A. Sanchez-Hernandez,[90] S. Carrillo Moreno,[91] C. Oropeza Barrera,[91] F. Vazquez Valencia,[91] J. Eysermans,[92] I. Pedraza,[92] H. A. Salazar Ibarguen,[92] C. Uribe Estrada,[92] A. Morelos Pineda,[93] D. Krofcheck,[94] S. Bheesette,[95] P. H. Butler,[95] A. Ahmad,[96] M. Ahmad,[96] M. I. Asghar,[96] Q. Hassan,[96] H. R. Hoorani,[96] A. Saddique,[96] M. A. Shah,[96] M. Shoaib,[96] M. Waqas,[96] H. Bialkowska,[97] M. Bluj,[97] B. Boimska,[97] T. Frueboes,[97] M. Górski,[97] M. Kazana,[97] M. Szleper,[97] P. Traczyk,[97] P. Zalewski,[97] K. Bunkowski,[98] A. Byszuk,[98,ii] K. Doroba,[98] A. Kalinowski,[98] M. Konecki,[98] J. Krolikowski,[98] M. Misiura,[98] M. Olszewski,[98] A. Pyskir,[98] M. Walczak,[98] M. Araujo,[99] P. Bargassa,[99] C. Beirão Da Cruz E Silva,[99] A. Di Francesco,[99] P. Faccioli,[99] B. Galinhas,[99] M. Gallinaro,[99] J. Hollar,[99] N. Leonardo,[99] M. V. Nemallapudi,[99] J. Seixas,[99] G. Strong,[99] O. Toldaiev,[99] D. Vadruccio,[99] J. Varela,[99] S. Afanasiev,[100] P. Bunin,[100] M. Gavrilenko,[100] I. Golutvin,[100] I. Gorbunov,[100] A. Kamenev,[100] V. Karjavine,[100] A. Lanev,[100] A. Malakhov,[100] V. Matveev,[100,jj,kk] P. Moisenz,[100] V. Palichik,[100] V. Perelygin,[100] S. Shmatov,[100] S. Shulha,[100] N. Skatchkov,[100] V. Smirnov,[100] N. Voytishin,[100] A. Zarubin,[100] V. Golovtsov,[101] Y. Ivanov,[101] V. Kim,[101,ll] E. Kuznetsova,[101,mm] P. Levchenko,[101] V. Murzin,[101] V. Oreshkin,[101] I. Smirnov,[101] D. Sosnov,[101] V. Sulimov,[101] L. Uvarov,[101] S. Vavilov,[101] A. Vorobyev,[101] Yu. Andreev,[102] A. Dermenev,[102] S. Gninenko,[102] N. Golubev,[102] A. Karneyeu,[102] M. Kirsanov,[102] N. Krasnikov,[102] A. Pashenkov,[102] D. Tlisov,[102] A. Toropin,[102] V. Epshteyn,[103] V. Gavrilov,[103] N. Lychkovskaya,[103] V. Popov,[103] I. Pozdnyakov,[103] G. Safronov,[103] A. Spiridonov,[103] A. Stepennov,[103] V. Stolin,[103] M. Toms,[103] E. Vlasov,[103] A. Zhokin,[103] T. Aushev,[104] M. Chadeeva,[105,nn] P. Parygin,[105] D. Philippov,[105] S. Polikarpov,[105,nn] E. Popova,[105] V. Rusinov,[105] V. Andreev,[106] M. Azarkin,[106] I. Dremin,[106,kk] M. Kirakosyan,[106] S. V. Rusakov,[106] A. Terkulov,[106] A. Baskakov,[107] A. Belyaev,[107] E. Boos,[107] V. Bunichev,[107] M. Dubinin,[107,oo] L. Dudko,[107] A. Ershov,[107] A. Gribushin,[107] V. Klyukhin,[107] O. Kodolova,[107] I. Lokhtin,[107] I. Miagkov,[107] S. Obraztsov,[107] S. Petrushanko,[107] V. Savrin,[107] A. Barnyakov,[108,pp] V. Blinov,[108,pp] T. Dimova,[108,pp] L. Kardapoltsev,[108,pp] Y. Skovpen,[108,pp] I. Azhgirey,[109] I. Bayshev,[109] S. Bitioukov,[109] D. Elumakhov,[109] A. Godizov,[109] V. Kachanov,[109] A. Kalinin,[109] D. Konstantinov,[109] P. Mandrik,[109] V. Petrov,[109] R. Ryutin,[109] S. Slabospitskii,[109] A. Sobol,[109] S. Troshin,[109] N. Tyurin,[109] A. Uzunian,[109] A. Volkov,[109] A. Babaev,[110] S. Baidali,[110] V. Okhotnikov,[110] P. Adzic,[111,qq] P. Cirkovic,[111] D. Devetak,[111] M. Dordevic,[111] J. Milosevic,[111] J. Alcaraz Maestre,[112] A. Álvarez Fernández,[112] I. Bachiller,[112] M. Barrio Luna,[112]

J. A. Brochero Cifuentes,[112] M. Cerrada,[112] N. Colino,[112] B. De La Cruz,[112] A. Delgado Peris,[112] C. Fernandez Bedoya,[112] J. P. Fernández Ramos,[112] J. Flix,[112] M. C. Fouz,[112] O. Gonzalez Lopez,[112] S. Goy Lopez,[112] J. M. Hernandez,[112] M. I. Josa,[112] D. Moran,[112] A. Pérez-Calero Yzquierdo,[112] J. Puerta Pelayo,[112] I. Redondo,[112] L. Romero,[112] M. S. Soares,[112] A. Triossi,[112] C. Albajar,[113] J. F. de Trocóniz,[113] J. Cuevas,[114] C. Erice,[114] J. Fernandez Menendez,[114] S. Folgueras,[114] I. Gonzalez Caballero,[114] J. R. González Fernández,[114] E. Palencia Cortezon,[114] V. Rodríguez Bouza,[114] S. Sanchez Cruz,[114] P. Vischia,[114] J. M. Vizan Garcia,[114] I. J. Cabrillo,[115] A. Calderon,[115] B. Chazin Quero,[115] J. Duarte Campderros,[115] M. Fernandez,[115] P. J. Fernández Manteca,[115] A. García Alonso,[115] J. Garcia-Ferrero,[115] G. Gomez,[115] A. Lopez Virto,[115] J. Marco,[115] C. Martinez Rivero,[115] P. Martinez Ruiz del Arbol,[115] F. Matorras,[115] J. Piedra Gomez,[115] C. Prieels,[115] T. Rodrigo,[115] A. Ruiz-Jimeno,[115] L. Scodellaro,[115] N. Trevisani,[115] I. Vila,[115] R. Vilar Cortabitarte,[115] N. Wickramage,[116] D. Abbaneo,[117] B. Akgun,[117] E. Auffray,[117] G. Auzinger,[117] P. Baillon,[117] A. H. Ball,[117] D. Barney,[117] J. Bendavid,[117] M. Bianco,[117] A. Bocci,[117] C. Botta,[117] E. Brondolin,[117] T. Camporesi,[117] M. Cepeda,[117] G. Cerminara,[117] E. Chapon,[117] Y. Chen,[117] G. Cucciati,[117] D. d'Enterria,[117] A. Dabrowski,[117] N. Daci,[117] V. Daponte,[117] A. David,[117] A. De Roeck,[117] N. Deelen,[117] M. Dobson,[117] M. Dünser,[117] N. Dupont,[117] A. Elliott-Peisert,[117] P. Everaerts,[117] F. Fallavollita,[117,rr] D. Fasanella,[117] G. Franzoni,[117] J. Fulcher,[117] W. Funk,[117] D. Gigi,[117] A. Gilbert,[117] K. Gill,[117] F. Glege,[117] M. Guilbaud,[117] D. Gulhan,[117] J. Hegeman,[117] C. Heidegger,[117] V. Innocente,[117] A. Jafari,[117] P. Janot,[117] O. Karacheban,[117,t] J. Kieseler,[117] A. Kornmayer,[117] M. Krammer,[117,b] C. Lange,[117] P. Lecoq,[117] C. Lourenço,[117] L. Malgeri,[117] M. Mannelli,[117] F. Meijers,[117] J. A. Merlin,[117] S. Mersi,[117] E. Meschi,[117] P. Milenovic,[117,ss] F. Moortgat,[117] M. Mulders,[117] J. Ngadiuba,[117] S. Nourbakhsh,[117] S. Orfanelli,[117] L. Orsini,[117] F. Pantaleo,[117,q] L. Pape,[117] E. Perez,[117] M. Peruzzi,[117] A. Petrilli,[117] G. Petrucciani,[117] A. Pfeiffer,[117] M. Pierini,[117] F. M. Pitters,[117] D. Rabady,[117] A. Racz,[117] T. Reis,[117] G. Rolandi,[117,tt] M. Rovere,[117] H. Sakulin,[117] C. Schäfer,[117] C. Schwick,[117] M. Seidel,[117] M. Selvaggi,[117] A. Sharma,[117] P. Silva,[117] P. Sphicas,[117,uu] A. Stakia,[117] J. Steggemann,[117] M. Tosi,[117] D. Treille,[117] A. Tsirou,[117] V. Veckalns,[117,vv] M. Verzetti,[117] W. D. Zeuner,[117] L. Caminada,[118,ww] K. Deiters,[118] W. Erdmann,[118] R. Horisberger,[118] Q. Ingram,[118] H. C. Kaestli,[118] D. Kotlinski,[118] U. Langenegger,[118] T. Rohe,[118] S. A. Wiederkehr,[118] M. Backhaus,[119] L. Bäni,[119] P. Berger,[119] N. Chernyavskaya,[119] G. Dissertori,[119] M. Dittmar,[119] M. Donegà,[119] C. Dorfer,[119] T. A. Gómez Espinosa,[119] C. Grab,[119] D. Hits,[119] T. Klijnsma,[119] W. Lustermann,[119] R. A. Manzoni,[119] M. Marionneau,[119] M. T. Meinhard,[119] F. Micheli,[119] P. Musella,[119] F. Nessi-Tedaldi,[119] J. Pata,[119] F. Pauss,[119] G. Perrin,[119] L. Perrozzi,[119] S. Pigazzini,[119] M. Quittnat,[119] C. Reissel,[119] D. Ruini,[119] D. A. Sanz Becerra,[119] M. Schönenberger,[119] L. Shchutska,[119] V. R. Tavolaro,[119] K. Theofilatos,[119] M. L. Vesterbacka Olsson,[119] R. Wallny,[119] D. H. Zhu,[119] T. K. Aarrestad,[120] C. Amsler,[120,xx] D. Brzhechko,[120] M. F. Canelli,[120] A. De Cosa,[120] R. Del Burgo,[120] S. Donato,[120] C. Galloni,[120] T. Hreus,[120] B. Kilminster,[120] S. Leontsinis,[120] I. Neutelings,[120] G. Rauco,[120] P. Robmann,[120] D. Salerno,[120] K. Schweiger,[120] C. Seitz,[120] Y. Takahashi,[120] A. Zucchetta,[120] Y. H. Chang,[121] K. y. Cheng,[121] T. H. Doan,[121] R. Khurana,[121] C. M. Kuo,[121] W. Lin,[121] A. Pozdnyakov,[121] S. S. Yu,[121] P. Chang,[122] Y. Chao,[122] K. F. Chen,[122] P. H. Chen,[122] W.-S. Hou,[122] Arun Kumar,[122] Y. F. Liu,[122] R.-S. Lu,[122] E. Paganis,[122] A. Psallidas,[122] A. Steen,[122] B. Asavapibhop,[123] N. Srimanobhas,[123] N. Suwonjandee,[123] A. Bat,[124] F. Boran,[124] S. Cerci,[124,yy] S. Damarseckin,[124] Z. S. Demiroglu,[124] F. Dolek,[124] C. Dozen,[124] I. Dumanoglu,[124] S. Girgis,[124] G. Gokbulut,[124] Y. Guler,[124] E. Gurpinar,[124] I. Hos,[124,zz] C. Isik,[124] E. E. Kangal,[124,aaa] O. Kara,[124] A. Kayis Topaksu,[124] U. Kiminsu,[124] M. Oglakci,[124] G. Onengut,[124] K. Ozdemir,[124,bbb] S. Ozturk,[124,ccc] D. Sunar Cerci,[124,yy] B. Tali,[124,yy] U. G. Tok,[124] S. Turkcapar,[124] I. S. Zorbakir,[124] C. Zorbilmez,[124] B. Isildak,[125,ddd] G. Karapinar,[125,eee] M. Yalvac,[125] M. Zeyrek,[125] I. O. Atakisi,[126] E. Gülmez,[126] M. Kaya,[126,fff] O. Kaya,[126,ggg] S. Ozkorucuklu,[126,hhh] S. Tekten,[126] E. A. Yetkin,[126,iii] M. N. Agaras,[127] A. Cakir,[127] K. Cankocak,[127] Y. Komurcu,[127] S. Sen,[127,jjj] B. Grynyov,[128] L. Levchuk,[129] F. Ball,[130] L. Beck,[130] J. J. Brooke,[130] D. Burns,[130] E. Clement,[130] D. Cussans,[130] O. Davignon,[130] H. Flacher,[130] J. Goldstein,[130] G. P. Heath,[130] H. F. Heath,[130] L. Kreczko,[130] D. M. Newbold,[130,kkk] S. Paramesvaran,[130] B. Penning,[130] T. Sakuma,[130] D. Smith,[130] V. J. Smith,[130] J. Taylor,[130] A. Titterton,[130] K. W. Bell,[131] A. Belyaev,[131,lll] C. Brew,[131] R. M. Brown,[131] D. Cieri,[131] D. J. A. Cockerill,[131] J. A. Coughlan,[131] K. Harder,[131] S. Harper,[131] J. Linacre,[131] E. Olaiya,[131] D. Petyt,[131] C. H. Shepherd-Themistocleous,[131] A. Thea,[131] I. R. Tomalin,[131] T. Williams,[131] W. J. Womersley,[131] R. Bainbridge,[132] P. Bloch,[132] J. Borg,[132] S. Breeze,[132] O. Buchmuller,[132] A. Bundock,[132] D. Colling,[132] P. Dauncey,[132] G. Davies,[132] M. Della Negra,[132] R. Di Maria,[132] Y. Haddad,[132] G. Hall,[132] G. Iles,[132] T. James,[132] M. Komm,[132] C. Laner,[132] L. Lyons,[132] A.-M. Magnan,[132] S. Malik,[132] A. Martelli,[132] J. Nash,[132,mmm] A. Nikitenko,[132,h] V. Palladino,[132] M. Pesaresi,[132] D. M. Raymond,[132] A. Richards,[132] A. Rose,[132] E. Scott,[132] C. Seez,[132] A. Shtipliyski,[132] G. Singh,[132] M. Stoye,[132] T. Strebler,[132] S. Summers,[132] A. Tapper,[132] K. Uchida,[132] T. Virdee,[132,q] N. Wardle,[132]

D. Winterbottom,[132] J. Wright,[132] S. C. Zenz,[132] J. E. Cole,[133] P. R. Hobson,[133] A. Khan,[133] P. Kyberd,[133] C. K. Mackay,[133] A. Morton,[133] I. D. Reid,[133] L. Teodorescu,[133] S. Zahid,[133] K. Call,[134] J. Dittmann,[134] K. Hatakeyama,[134] H. Liu,[134] C. Madrid,[134] B. Mcmaster,[134] N. Pastika,[134] C. Smith,[134] R. Bartek,[135] A. Dominguez,[135] A. Buccilli,[136] S. I. Cooper,[136] C. Henderson,[136] P. Rumerio,[136] C. West,[136] D. Arcaro,[137] T. Bose,[137] D. Gastler,[137] D. Pinna,[137] D. Rankin,[137] C. Richardson,[137] J. Rohlf,[137] L. Sulak,[137] D. Zou,[137] G. Benelli,[138] X. Coubez,[138] D. Cutts,[138] M. Hadley,[138] J. Hakala,[138] U. Heintz,[138] J. M. Hogan,[138,nnn] K. H. M. Kwok,[138] E. Laird,[138] G. Landsberg,[138] J. Lee,[138] Z. Mao,[138] M. Narain,[138] S. Sagir,[138,ooo] R. Syarif,[138] E. Usai,[138] D. Yu,[138] R. Band,[139] C. Brainerd,[139] R. Breedon,[139] D. Burns,[139] M. Calderon De La Barca Sanchez,[139] M. Chertok,[139] J. Conway,[139] R. Conway,[139] P. T. Cox,[139] R. Erbacher,[139] C. Flores,[139] G. Funk,[139] W. Ko,[139] O. Kukral,[139] R. Lander,[139] M. Mulhearn,[139] D. Pellett,[139] J. Pilot,[139] S. Shalhout,[139] M. Shi,[139] D. Stolp,[139] D. Taylor,[139] K. Tos,[139] M. Tripathi,[139] Z. Wang,[139] F. Zhang,[139] M. Bachtis,[140] C. Bravo,[140] R. Cousins,[140] A. Dasgupta,[140] A. Florent,[140] J. Hauser,[140] M. Ignatenko,[140] N. Mccoll,[140] S. Regnard,[140] D. Saltzberg,[140] C. Schnaible,[140] V. Valuev,[140] E. Bouvier,[141] K. Burt,[141] R. Clare,[141] J. W. Gary,[141] S. M. A. Ghiasi Shirazi,[141] G. Hanson,[141] G. Karapostoli,[141] E. Kennedy,[141] F. Lacroix,[141] O. R. Long,[141] M. Olmedo Negrete,[141] M. I. Paneva,[141] W. Si,[141] L. Wang,[141] H. Wei,[141] S. Wimpenny,[141] B. R. Yates,[141] J. G. Branson,[142] P. Chang,[142] S. Cittolin,[142] M. Derdzinski,[142] R. Gerosa,[142] D. Gilbert,[142] B. Hashemi,[142] A. Holzner,[142] D. Klein,[142] G. Kole,[142] V. Krutelyov,[142] J. Letts,[142] M. Masciovecchio,[142] D. Olivito,[142] S. Padhi,[142] M. Pieri,[142] M. Sani,[142] V. Sharma,[142] S. Simon,[142] M. Tadel,[142] A. Vartak,[142] S. Wasserbaech,[142,ppp] J. Wood,[142] F. Würthwein,[142] A. Yagil,[142] G. Zevi Della Porta,[142] N. Amin,[143] R. Bhandari,[143] J. Bradmiller-Feld,[143] C. Campagnari,[143] M. Citron,[143] A. Dishaw,[143] V. Dutta,[143] M. Franco Sevilla,[143] L. Gouskos,[143] R. Heller,[143] J. Incandela,[143] A. Ovcharova,[143] H. Qu,[143] J. Richman,[143] D. Stuart,[143] I. Suarez,[143] S. Wang,[143] J. Yoo,[143] D. Anderson,[144] A. Bornheim,[144] J. M. Lawhorn,[144] H. B. Newman,[144] T. Q. Nguyen,[144] M. Spiropulu,[144] J. R. Vlimant,[144] R. Wilkinson,[144] S. Xie,[144] Z. Zhang,[144] R. Y. Zhu,[144] M. B. Andrews,[145] T. Ferguson,[145] T. Mudholkar,[145] M. Paulini,[145] M. Sun,[145] I. Vorobiev,[145] M. Weinberg,[145] J. P. Cumalat,[146] W. T. Ford,[146] F. Jensen,[146] A. Johnson,[146] M. Krohn,[146] E. MacDonald,[146] T. Mulholland,[146] R. Patel,[146] A. Perloff,[146] K. Stenson,[146] K. A. Ulmer,[146] S. R. Wagner,[146] J. Alexander,[147] J. Chaves,[147] Y. Cheng,[147] J. Chu,[147] A. Datta,[147] K. Mcdermott,[147] N. Mirman,[147] J. R. Patterson,[147] D. Quach,[147] A. Rinkevicius,[147] A. Ryd,[147] L. Skinnari,[147] L. Soffi,[147] S. M. Tan,[147] Z. Tao,[147] J. Thom,[147] J. Tucker,[147] P. Wittich,[147] M. Zientek,[147] S. Abdullin,[148] M. Albrow,[148] M. Alyari,[148] G. Apollinari,[148] A. Apresyan,[148] A. Apyan,[148] S. Banerjee,[148] L. A. T. Bauerdick,[148] A. Beretvas,[148] J. Berryhill,[148] P. C. Bhat,[148] K. Burkett,[148] J. N. Butler,[148] A. Canepa,[148] G. B. Cerati,[148] H. W. K. Cheung,[148] F. Chlebana,[148] M. Cremonesi,[148] J. Duarte,[148] V. D. Elvira,[148] J. Freeman,[148] Z. Gecse,[148] E. Gottschalk,[148] L. Gray,[148] D. Green,[148] S. Grünendahl,[148] O. Gutsche,[148] J. Hanlon,[148] R. M. Harris,[148] S. Hasegawa,[148] J. Hirschauer,[148] Z. Hu,[148] B. Jayatilaka,[148] S. Jindariani,[148] M. Johnson,[148] U. Joshi,[148] B. Klima,[148] M. J. Kortelainen,[148] B. Kreis,[148] S. Lammel,[148] D. Lincoln,[148] R. Lipton,[148] M. Liu,[148] T. Liu,[148] J. Lykken,[148] K. Maeshima,[148] J. M. Marraffino,[148] D. Mason,[148] P. McBride,[148] P. Merkel,[148] S. Mrenna,[148] S. Nahn,[148] V. O'Dell,[148] K. Pedro,[148] C. Pena,[148] O. Prokofyev,[148] G. Rakness,[148] L. Ristori,[148] A. Savoy-Navarro,[148,qqq] B. Schneider,[148] E. Sexton-Kennedy,[148] A. Soha,[148] W. J. Spalding,[148] L. Spiegel,[148] S. Stoynev,[148] J. Strait,[148] N. Strobbe,[148] L. Taylor,[148] S. Tkaczyk,[148] N. V. Tran,[148] L. Uplegger,[148] E. W. Vaandering,[148] C. Vernieri,[148] M. Verzocchi,[148] R. Vidal,[148] M. Wang,[148] H. A. Weber,[148] A. Whitbeck,[148] D. Acosta,[149] P. Avery,[149] P. Bortignon,[149] D. Bourilkov,[149] A. Brinkerhoff,[149] L. Cadamuro,[149] A. Carnes,[149] M. Carver,[149] D. Curry,[149] R. D. Field,[149] S. V. Gleyzer,[149] B. M. Joshi,[149] J. Konigsberg,[149] A. Korytov,[149] K. H. Lo,[149] P. Ma,[149] K. Matchev,[149] H. Mei,[149] G. Mitselmakher,[149] D. Rosenzweig,[149] K. Shi,[149] D. Sperka,[149] J. Wang,[149] S. Wang,[149] X. Zuo,[149] Y. R. Joshi,[150] S. Linn,[150] A. Ackert,[151] T. Adams,[151] A. Askew,[151] S. Hagopian,[151] V. Hagopian,[151] K. F. Johnson,[151] T. Kolberg,[151] G. Martinez,[151] T. Perry,[151] H. Prosper,[151] A. Saha,[151] C. Schiber,[151] R. Yohay,[151] M. M. Baarmand,[152] V. Bhopatkar,[152] S. Colafranceschi,[152] M. Hohlmann,[152] D. Noonan,[152] M. Rahmani,[152] T. Roy,[152] F. Yumiceva,[152] M. R. Adams,[153] L. Apanasevich,[153] D. Berry,[153] R. R. Betts,[153] R. Cavanaugh,[153] X. Chen,[153] S. Dittmer,[153] O. Evdokimov,[153] C. E. Gerber,[153] D. A. Hangal,[153] D. J. Hofman,[153] K. Jung,[153] J. Kamin,[153] C. Mills,[153] I. D. Sandoval Gonzalez,[153] M. B. Tonjes,[153] H. Trauger,[153] N. Varelas,[153] H. Wang,[153] X. Wang,[153] Z. Wu,[153] J. Zhang,[153] M. Alhusseini,[154] B. Bilki,[154,rrr] W. Clarida,[154] K. Dilsiz,[154,sss] S. Durgut,[154] R. P. Gandrajula,[154] M. Haytmyradov,[154] V. Khristenko,[154] J.-P. Merlo,[154] A. Mestvirishvili,[154] A. Moeller,[154] J. Nachtman,[154] H. Ogul,[154,ttt] Y. Onel,[154] F. Ozok,[154,uuu] A. Penzo,[154] C. Snyder,[154] E. Tiras,[154] J. Wetzel,[154] B. Blumenfeld,[155] A. Cocoros,[155] N. Eminizer,[155] D. Fehling,[155] L. Feng,[155] A. V. Gritsan,[155] W. T. Hung,[155] P. Maksimovic,[155] J. Roskes,[155] U. Sarica,[155] M. Swartz,[155] M. Xiao,[155] C. You,[155] A. Al-bataineh,[156] P. Baringer,[156]

A. Bean,[156] S. Boren,[156] J. Bowen,[156] A. Bylinkin,[156] J. Castle,[156] S. Khalil,[156] A. Kropivnitskaya,[156] D. Majumder,[156] W. Mcbrayer,[156] M. Murray,[156] C. Rogan,[156] S. Sanders,[156] E. Schmitz,[156] J. D. Tapia Takaki,[156] Q. Wang,[156] S. Duric,[157] A. Ivanov,[157] K. Kaadze,[157] D. Kim,[157] Y. Maravin,[157] D. R. Mendis,[157] T. Mitchell,[157] A. Modak,[157] A. Mohammadi,[157] L. K. Saini,[157] N. Skhirtladze,[157] F. Rebassoo,[158] D. Wright,[158] A. Baden,[159] O. Baron,[159] A. Belloni,[159] S. C. Eno,[159] Y. Feng,[159] C. Ferraioli,[159] N. J. Hadley,[159] S. Jabeen,[159] G. Y. Jeng,[159] R. G. Kellogg,[159] J. Kunkle,[159] A. C. Mignerey,[159] S. Nabili,[159] F. Ricci-Tam,[159] Y. H. Shin,[159] A. Skuja,[159] S. C. Tonwar,[159] K. Wong,[159] D. Abercrombie,[160] B. Allen,[160] V. Azzolini,[160] A. Baty,[160] G. Bauer,[160] R. Bi,[160] S. Brandt,[160] W. Busza,[160] I. A. Cali,[160] M. D'Alfonso,[160] Z. Demiragli,[160] G. Gomez Ceballos,[160] M. Goncharov,[160] P. Harris,[160] D. Hsu,[160] M. Hu,[160] Y. Iiyama,[160] G. M. Innocenti,[160] M. Klute,[160] D. Kovalskyi,[160] Y.-J. Lee,[160] P. D. Luckey,[160] B. Maier,[160] A. C. Marini,[160] C. Mcginn,[160] C. Mironov,[160] S. Narayanan,[160] X. Niu,[160] C. Paus,[160] C. Roland,[160] G. Roland,[160] G. S. F. Stephans,[160] K. Sumorok,[160] K. Tatar,[160] D. Velicanu,[160] J. Wang,[160] T. W. Wang,[160] B. Wyslouch,[160] S. Zhaozhong,[160] A. C. Benvenuti,[161,a] R. M. Chatterjee,[161] A. Evans,[161] P. Hansen,[161] J. Hiltbrand,[161] Sh. Jain,[161] S. Kalafut,[161] Y. Kubota,[161] Z. Lesko,[161] J. Mans,[161] N. Ruckstuhl,[161] R. Rusack,[161] M. A. Wadud,[161] J. G. Acosta,[162] S. Oliveros,[162] E. Avdeeva,[163] K. Bloom,[163] D. R. Claes,[163] C. Fangmeier,[163] F. Golf,[163] R. Gonzalez Suarez,[163] R. Kamalieddin,[163] I. Kravchenko,[163] J. Monroy,[163] J. E. Siado,[163] G. R. Snow,[163] B. Stieger,[163] A. Godshalk,[164] C. Harrington,[164] I. Iashvili,[164] A. Kharchilava,[164] C. Mclean,[164] D. Nguyen,[164] A. Parker,[164] S. Rappoccio,[164] B. Roozbahani,[164] E. Barberis,[165] C. Freer,[165] A. Hortiangtham,[165] D. M. Morse,[165] T. Orimoto,[165] R. Teixeira De Lima,[165] T. Wamorkar,[165] B. Wang,[165] A. Wisecarver,[165] D. Wood,[165] S. Bhattacharya,[166] O. Charaf,[166] K. A. Hahn,[166] N. Mucia,[166] N. Odell,[166] M. H. Schmitt,[166] K. Sung,[166] M. Trovato,[166] M. Velasco,[166] R. Bucci,[167] N. Dev,[167] M. Hildreth,[167] K. Hurtado Anampa,[167] C. Jessop,[167] D. J. Karmgard,[167] N. Kellams,[167] K. Lannon,[167] W. Li,[167] N. Loukas,[167] N. Marinelli,[167] F. Meng,[167] C. Mueller,[167] Y. Musienko,[167,jj] M. Planer,[167] A. Reinsvold,[167] R. Ruchti,[167] P. Siddireddy,[167] G. Smith,[167] S. Taroni,[167] M. Wayne,[167] A. Wightman,[167] M. Wolf,[167] A. Woodard,[167] J. Alimena,[168] L. Antonelli,[168] B. Bylsma,[168] L. S. Durkin,[168] S. Flowers,[168] B. Francis,[168] A. Hart,[168] C. Hill,[168] W. Ji,[168] T. Y. Ling,[168] W. Luo,[168] B. L. Winer,[168] S. Cooperstein,[169] P. Elmer,[169] J. Hardenbrook,[169] S. Higginbotham,[169] A. Kalogeropoulos,[169] D. Lange,[169] M. T. Lucchini,[169] J. Luo,[169] D. Marlow,[169] K. Mei,[169] I. Ojalvo,[169] J. Olsen,[169] C. Palmer,[169] P. Piroué,[169] J. Salfeld-Nebgen,[169] D. Stickland,[169] C. Tully,[169] S. Malik,[170] S. Norberg,[170] A. Barker,[171] V. E. Barnes,[171] S. Das,[171] L. Gutay,[171] M. Jones,[171] A. W. Jung,[171] A. Khatiwada,[171] B. Mahakud,[171] D. H. Miller,[171] N. Neumeister,[171] C. C. Peng,[171] S. Piperov,[171] H. Qiu,[171] J. F. Schulte,[171] J. Sun,[171] F. Wang,[171] R. Xiao,[171] W. Xie,[171] T. Cheng,[172] J. Dolen,[172] N. Parashar,[172] Z. Chen,[173] K. M. Ecklund,[173] S. Freed,[173] F. J. M. Geurts,[173] M. Kilpatrick,[173] W. Li,[173] B. P. Padley,[173] R. Redjimi,[173] J. Roberts,[173] J. Rorie,[173] W. Shi,[173] Z. Tu,[173] J. Zabel,[173] A. Zhang,[173] A. Bodek,[174] P. de Barbaro,[174] R. Demina,[174] Y. t. Duh,[174] J. L. Dulemba,[174] C. Fallon,[174] T. Ferbel,[174] M. Galanti,[174] A. Garcia-Bellido,[174] J. Han,[174] O. Hindrichs,[174] A. Khukhunaishvili,[174] P. Tan,[174] R. Taus,[174] A. Agapitos,[175] J. P. Chou,[175] Y. Gershtein,[175] E. Halkiadakis,[175] M. Heindl,[175] E. Hughes,[175] S. Kaplan,[175] R. Kunnawalkam Elayavalli,[175] S. Kyriacou,[175] A. Lath,[175] R. Montalvo,[175] K. Nash,[175] M. Osherson,[175] H. Saka,[175] S. Salur,[175] S. Schnetzer,[175] D. Sheffield,[175] S. Somalwar,[175] R. Stone,[175] S. Thomas,[175] P. Thomassen,[175] M. Walker,[175] A. G. Delannoy,[176] J. Heideman,[176] G. Riley,[176] S. Spanier,[176] O. Bouhali,[177,vvv] A. Celik,[177] M. Dalchenko,[177] M. De Mattia,[177] A. Delgado,[177] S. Dildick,[177] R. Eusebi,[177] J. Gilmore,[177] T. Huang,[177] T. Kamon,[177,www] S. Luo,[177] R. Mueller,[177] D. Overton,[177] L. Perniè,[177] D. Rathjens,[177] A. Safonov,[177] N. Akchurin,[178] J. Damgov,[178] F. De Guio,[178] P. R. Dudero,[178] S. Kunori,[178] K. Lamichhane,[178] S. W. Lee,[178] T. Mengke,[178] S. Muthumuni,[178] T. Peltola,[178] S. Undleeb,[178] I. Volobouev,[178] Z. Wang,[178] S. Greene,[179] A. Gurrola,[179] R. Janjam,[179] W. Johns,[179] C. Maguire,[179] A. Melo,[179] H. Ni,[179] K. Padeken,[179] J. D. Ruiz Alvarez,[179] P. Sheldon,[179] S. Tuo,[179] J. Velkovska,[179] M. Verweij,[179] Q. Xu,[179] M. W. Arenton,[180] P. Barria,[180] B. Cox,[180] R. Hirosky,[180] M. Joyce,[180] A. Ledovskoy,[180] H. Li,[180] C. Neu,[180] T. Sinthuprasith,[180] Y. Wang,[180] E. Wolfe,[180] F. Xia,[180] R. Harr,[181] P. E. Karchin,[181] N. Poudyal,[181] J. Sturdy,[181] P. Thapa,[181] S. Zaleski,[181] M. Brodski,[182] J. Buchanan,[182] C. Caillol,[182] D. Carlsmith,[182] S. Dasu,[182] L. Dodd,[182] B. Gomber,[182] M. Grothe,[182] M. Herndon,[182] A. Hervé,[182] U. Hussain,[182] P. Klabbers,[182] A. Lanaro,[182] K. Long,[182] R. Loveless,[182] T. Ruggles,[182] A. Savin,[182] V. Sharma,[182] N. Smith,[182] W. H. Smith,[182] and N. Woods[182]

(CMS Collaboration)

[1]Yerevan Physics Institute, Yerevan, Armenia
[2]Institut für Hochenergiephysik, Wien, Austria
[3]Institute for Nuclear Problems, Minsk, Belarus
[4]Universiteit Antwerpen, Antwerpen, Belgium
[5]Vrije Universiteit Brussel, Brussel, Belgium
[6]Université Libre de Bruxelles, Bruxelles, Belgium
[7]Ghent University, Ghent, Belgium
[8]Université Catholique de Louvain, Louvain-la-Neuve, Belgium
[9]Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil
[10]Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil
[11]Universidade Estadual Paulista, Universidade Federal do ABC, São Paulo, Brazil
[11a]Universidade Estadual Paulista, São Paulo, Brazil
[11b]Universidade Federal do ABC, São Paulo, Brazil
[12]Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria
[13]University of Sofia, Sofia, Bulgaria
[14]Beihang University, Beijing, China
[15]Institute of High Energy Physics, Beijing, China
[16]State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China
[17]Tsinghua University, Beijing, China
[18]Universidad de Los Andes, Bogota, Colombia
[19]University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia
[20]University of Split, Faculty of Science, Split, Croatia
[21]Institute Rudjer Boskovic, Zagreb, Croatia
[22]University of Cyprus, Nicosia, Cyprus
[23]Charles University, Prague, Czech Republic
[24]Escuela Politecnica Nacional, Quito, Ecuador
[25]Universidad San Francisco de Quito, Quito, Ecuador
[26]Academy of Scientific Research and Technology of the Arab Republic of Egypt,
Egyptian Network of High Energy Physics, Cairo, Egypt
[27]National Institute of Chemical Physics and Biophysics, Tallinn, Estonia
[28]Department of Physics, University of Helsinki, Helsinki, Finland
[29]Helsinki Institute of Physics, Helsinki, Finland
[30]Lappeenranta University of Technology, Lappeenranta, Finland
[31]IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
[32]Laboratoire Leprince-Ringuet, Ecole polytechnique, CNRS/IN2P3, Université Paris-Saclay, Palaiseau, France
[33]Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France
[34]Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules,
CNRS/IN2P3, Villeurbanne, France
[35]Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon,
Villeurbanne, France
[36]Georgian Technical University, Tbilisi, Georgia
[37]Tbilisi State University, Tbilisi, Georgia
[38]RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany
[39]RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany
[40]RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany
[41]Deutsches Elektronen-Synchrotron, Hamburg, Germany
[42]University of Hamburg, Hamburg, Germany
[43]Karlsruher Institut fuer Technology, Karlsruhe, Germany
[44]Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece
[45]National and Kapodistrian University of Athens, Athens, Greece
[46]National Technical University of Athens, Athens, Greece
[47]University of Ioánnina, Ioánnina, Greece
[48]MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary
[49]Wigner Research Centre for Physics, Budapest, Hungary
[50]Institute of Nuclear Research ATOMKI, Debrecen, Hungary
[51]Institute of Physics, University of Debrecen, Debrecen, Hungary
[52]Indian Institute of Science (IISc), Bangalore, India
[53]National Institute of Science Education and Research, HBNI, Bhubaneswar, India
[54]Panjab University, Chandigarh, India

[55]University of Delhi, Delhi, India
[56]Saha Institute of Nuclear Physics, HBNI, Kolkata, India
[57]Indian Institute of Technology Madras, Madras, India
[58]Bhabha Atomic Research Centre, Mumbai, India
[59]Tata Institute of Fundamental Research-A, Mumbai, India
[60]Tata Institute of Fundamental Research-B, Mumbai, India
[61]Indian Institute of Science Education and Research (IISER), Pune, India
[62]Institute for Research in Fundamental Sciences (IPM), Tehran, Iran
[63]University College Dublin, Dublin, Ireland
[64a]INFN Sezione di Bari, Bari, Italy
[64b]Università di Bari, Bari, Italy
[64c]Politecnico di Bari, Bari, Italy
[65a]INFN Sezione di Bologna, Bologna, Italy
[65b]Università di Bologna, Bologna, Italy
[66a]INFN Sezione di Catania, Catania, Italy
[66b]Università di Catania, Catania, Italy
[67a]INFN Sezione di Firenze, Firenze, Italy
[67b]Università di Firenze, Firenze, Italy
[68]INFN Laboratori Nazionali di Frascati, Frascati, Italy
[69a]INFN Sezione di Genova, Genova, Italy
[69b]Università di Genova, Genova, Italy
[70a]INFN Sezione di Milano-Bicocca, Milano, Italy
[70b]Università di Milano-Bicocca, Milano, Italy
[71a]INFN Sezione di Napoli, Napoli, Italy
[71b]Università di Napoli 'Federico II', Napoli, Italy
[71c]Università della Basilicata, Potenza, Italy
[71d]Università G. Marconi, Roma, Italy
[72a]INFN Sezione di Padova, Padova, Italy
[72b]Università di Padova, Padova, Italy
[72c]Università di Trento, Trento, Italy
[73a]INFN Sezione di Pavia, Pavia, Italy
[73b]Università di Pavia, Pavia, Italy
[74a]INFN Sezione di Perugia, Perugia, Italy
[74b]Università di Perugia, Perugia, Italy
[75a]INFN Sezione di Pisa, Pisa, Italy
[75b]Università di Pisa, Pisa, Italy
[75c]Scuola Normale Superiore di Pisa, Pisa, Italy
[76a]INFN Sezione di Roma, Rome, Italy
[76b]Sapienza Università di Roma, Rome, Italy
[77a]INFN Sezione di Torino, Torino, Italy
[77b]Università di Torino, Torino, Italy
[77c]Università del Piemonte Orientale, Novara, Italy
[78a]INFN Sezione di Trieste, Trieste, Italy
[78b]Università di Trieste, Trieste, Italy
[79]Kyungpook National University, Daegu, Korea
[80]Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea
[81]Hanyang University, Seoul, Korea
[82]Korea University, Seoul, Korea
[83]Sejong University, Seoul, Korea
[84]Seoul National University, Seoul, Korea
[85]University of Seoul, Seoul, Korea
[86]Sungkyunkwan University, Suwon, Korea
[87]Vilnius University, Vilnius, Lithuania
[88]National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia
[89]Universidad de Sonora (UNISON), Hermosillo, Mexico
[90]Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico
[91]Universidad Iberoamericana, Mexico City, Mexico
[92]Benemerita Universidad Autonoma de Puebla, Puebla, Mexico
[93]Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico
[94]University of Auckland, Auckland, New Zealand

[95]University of Canterbury, Christchurch, New Zealand
[96]National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan
[97]National Centre for Nuclear Research, Swierk, Poland
[98]Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland
[99]Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal
[100]Joint Institute for Nuclear Research, Dubna, Russia
[101]Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia
[102]Institute for Nuclear Research, Moscow, Russia
[103]Institute for Theoretical and Experimental Physics, Moscow, Russia
[104]Moscow Institute of Physics and Technology, Moscow, Russia
[105]National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia
[106]P.N. Lebedev Physical Institute, Moscow, Russia
[107]Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia
[108]Novosibirsk State University (NSU), Novosibirsk, Russia
[109]State Research Center of Russian Federation, Institute for High Energy Physics of NRC "Kurchatov Institute", Protvino, Russia
[110]National Research Tomsk Polytechnic University, Tomsk, Russia
[111]University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia
[112]Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain
[113]Universidad Autónoma de Madrid, Madrid, Spain
[114]Universidad de Oviedo, Oviedo, Spain
[115]Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain
[116]University of Ruhuna, Department of Physics, Matara, Sri Lanka
[117]CERN, European Organization for Nuclear Research, Geneva, Switzerland
[118]Paul Scherrer Institut, Villigen, Switzerland
[119]ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland
[120]Universität Zürich, Zurich, Switzerland
[121]National Central University, Chung-Li, Taiwan
[122]National Taiwan University (NTU), Taipei, Taiwan
[123]Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand
[124]Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey
[125]Middle East Technical University, Physics Department, Ankara, Turkey
[126]Bogazici University, Istanbul, Turkey
[127]Istanbul Technical University, Istanbul, Turkey
[128]Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkov, Ukraine
[129]National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine
[130]University of Bristol, Bristol, United Kingdom
[131]Rutherford Appleton Laboratory, Didcot, United Kingdom
[132]Imperial College, London, United Kingdom
[133]Brunel University, Uxbridge, United Kingdom
[134]Baylor University, Waco, Texas, USA
[135]Catholic University of America, Washington, DC, USA
[136]The University of Alabama, Tuscaloosa, Alabama, USA
[137]Boston University, Boston, Massachusetts, USA
[138]Brown University, Providence, Rhode Island, USA
[139]University of California, Davis, Davis, California, USA
[140]University of California, Los Angeles, California, USA
[141]University of California, Riverside, Riverside, California, USA
[142]University of California, San Diego, La Jolla, California, USA
[143]University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA
[144]California Institute of Technology, Pasadena, California, USA
[145]Carnegie Mellon University, Pittsburgh, Pennsylvania, USA
[146]University of Colorado Boulder, Boulder, Colorado, USA
[147]Cornell University, Ithaca, New York, USA
[148]Fermi National Accelerator Laboratory, Batavia, Illinois, USA
[149]University of Florida, Gainesville, Florida, USA
[150]Florida International University, Miami, Florida, USA
[151]Florida State University, Tallahassee, Florida, USA
[152]Florida Institute of Technology, Melbourne, Florida, USA
[153]University of Illinois at Chicago (UIC), Chicago, Illinois, USA

[154]The University of Iowa, Iowa City, Iowa, USA
[155]Johns Hopkins University, Baltimore, Maryland, USA
[156]The University of Kansas, Lawrence, Kansas, USA
[157]Kansas State University, Manhattan, Kansas, USA
[158]Lawrence Livermore National Laboratory, Livermore, California, USA
[159]University of Maryland, College Park, Maryland, USA
[160]Massachusetts Institute of Technology, Cambridge, Massachusetts, USA
[161]University of Minnesota, Minneapolis, Minnesota, USA
[162]University of Mississippi, Oxford, Mississippi, USA
[163]University of Nebraska-Lincoln, Lincoln, Nebraska, USA
[164]State University of New York at Buffalo, Buffalo, New York, USA
[165]Northeastern University, Boston, Massachusetts, USA
[166]Northwestern University, Evanston, Illinois, USA
[167]University of Notre Dame, Notre Dame, Indiana, USA
[168]The Ohio State University, Columbus, Ohio, USA
[169]Princeton University, Princeton, New Jersey, USA
[170]University of Puerto Rico, Mayaguez, Puerto Rico
[171]Purdue University, West Lafayette, Indiana, USA
[172]Purdue University Northwest, Hammond, Indiana, USA
[173]Rice University, Houston, Texas, USA
[174]University of Rochester, Rochester, New York, USA
[175]Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA
[176]University of Tennessee, Knoxville, Tennessee, USA
[177]Texas A&M University, College Station, Texas, USA
[178]Texas Tech University, Lubbock, Texas, USA
[179]Vanderbilt University, Nashville, Tennessee, USA
[180]University of Virginia, Charlottesville, Virginia, USA
[181]Wayne State University, Detroit, Michigan, USA
[182]University of Wisconsin - Madison, Madison, Wisconsin, USA

[a]Deceased.
[b]Also at Vienna University of Technology, Vienna, Austria.
[c]Also at IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France.
[d]Also at Universidade Estadual de Campinas, Campinas, Brazil.
[e]Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil.
[f]Also at Université Libre de Bruxelles, Bruxelles, Belgium.
[g]Also at University of Chinese Academy of Sciences, Beijing, China.
[h]Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.
[i]Also at Joint Institute for Nuclear Research, Dubna, Russia.
[j]Also at British University in Egypt, Cairo, Egypt.
[k]Also at Suez University, Suez, Egypt.
[l]Also at Zewail City of Science and Technology, Zewail, Egypt.
[m]Also at Department of Physics, King Abdulaziz University, Jeddah, Saudi Arabia.
[n]Also at Université de Haute Alsace, Mulhouse, France.
[o]Also at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia.
[p]Also at Tbilisi State University, Tbilisi, Georgia.
[q]Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
[r]Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany.
[s]Also at University of Hamburg, Hamburg, Germany.
[t]Also at Brandenburg University of Technology, Cottbus, Germany.
[u]Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary.
[v]Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
[w]Also at Institute of Physics, University of Debrecen, Debrecen, Hungary.
[x]Also at IIT Bhubaneswar, Bhubaneswar, India.
[y]Also at Institute of Physics, Bhubaneswar, India.
[z]Also at Shoolini University, Solan, India.
[aa]Also at University of Visva-Bharati, Santiniketan, India.
[bb]Also at Isfahan University of Technology, Isfahan, Iran.
[cc]Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
[dd]Also at Università degli Studi di Siena, Siena, Italy.

[ee]Also at Kyunghee University, Seoul, Korea.
[ff]Also at International Islamic University of Malaysia, Kuala Lumpur, Malaysia.
[gg]Also at Malaysian Nuclear Agency, MOSTI, Kajang, Malaysia.
[hh]Also at Consejo Nacional de Ciencia y Tecnología, Mexico city, Mexico.
[ii]Also at Warsaw University of Technology, Institute of Electronic Systems, Warsaw, Poland.
[jj]Also at Institute for Nuclear Research, Moscow, Russia.
[kk]Also at National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia.
[ll]Also at St. Petersburg State Polytechnical University, St. Petersburg, Russia.
[mm]Also at University of Florida, Gainesville, Florida, USA.
[nn]Also at P.N. Lebedev Physical Institute, Moscow, Russia.
[oo]Also at California Institute of Technology, Pasadena, California, USA.
[pp]Also at Budker Institute of Nuclear Physics, Novosibirsk, Russia.
[qq]Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
[rr]Also at INFN Sezione di Pavia, Università di Pavia, Pavia, Italy.
[ss]Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.
[tt]Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy.
[uu]Also at National and Kapodistrian University of Athens, Athens, Greece.
[vv]Also at Riga Technical University, Riga, Latvia.
[ww]Also at Universität Zürich, Zurich, Switzerland.
[xx]Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria.
[yy]Also at Adiyaman University, Adiyaman, Turkey.
[zz]Also at Istanbul Aydin University, Istanbul, Turkey.
[aaa]Also at Mersin University, Mersin, Turkey.
[bbb]Also at Piri Reis University, Istanbul, Turkey.
[ccc]Also at Gaziosmanpasa University, Tokat, Turkey.
[ddd]Also at Ozyegin University, Istanbul, Turkey.
[eee]Also at Izmir Institute of Technology, Izmir, Turkey.
[fff]Also at Marmara University, Istanbul, Turkey.
[ggg]Also at Kafkas University, Kars, Turkey.
[hhh]Also at Istanbul University, Faculty of Science, Istanbul, Turkey.
[iii]Also at Istanbul Bilgi University, Istanbul, Turkey.
[jjj]Also at Hacettepe University, Ankara, Turkey.
[kkk]Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.
[lll]Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
[mmm]Also at Monash University, Faculty of Science, Clayton, Australia.
[nnn]Also at Bethel University, St. Paul, Minnesota, USA.
[ooo]Also at Karamanoğlu Mehmetbey University, Karaman, Turkey.
[ppp]Also at Utah Valley University, Orem, Utah, USA.
[qqq]Also at Purdue University, West Lafayette, Indiana, USA.
[rrr]Also at Beykent University, Istanbul, Turkey.
[sss]Also at Bingol University, Bingol, Turkey.
[ttt]Also at Sinop University, Sinop, Turkey.
[uuu]Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
[vvv]Also at Texas A&M University at Qatar, Doha, Qatar.
[www]Also at Kyungpook National University, Daegu, Korea.